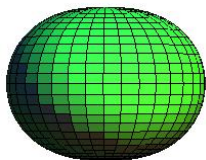


Effects of nuclear Deformation in *Heavy Ion Collisions.*

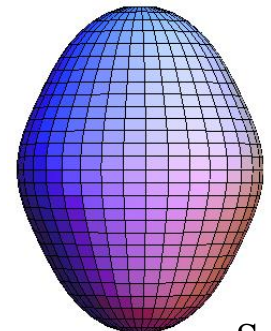
Peter FILIP

Institute of Physics, Slovak Academy of Sciences
Bratislava 845 11, Slovak Republic



Ge-70

Kent State University
April 10th, 2009



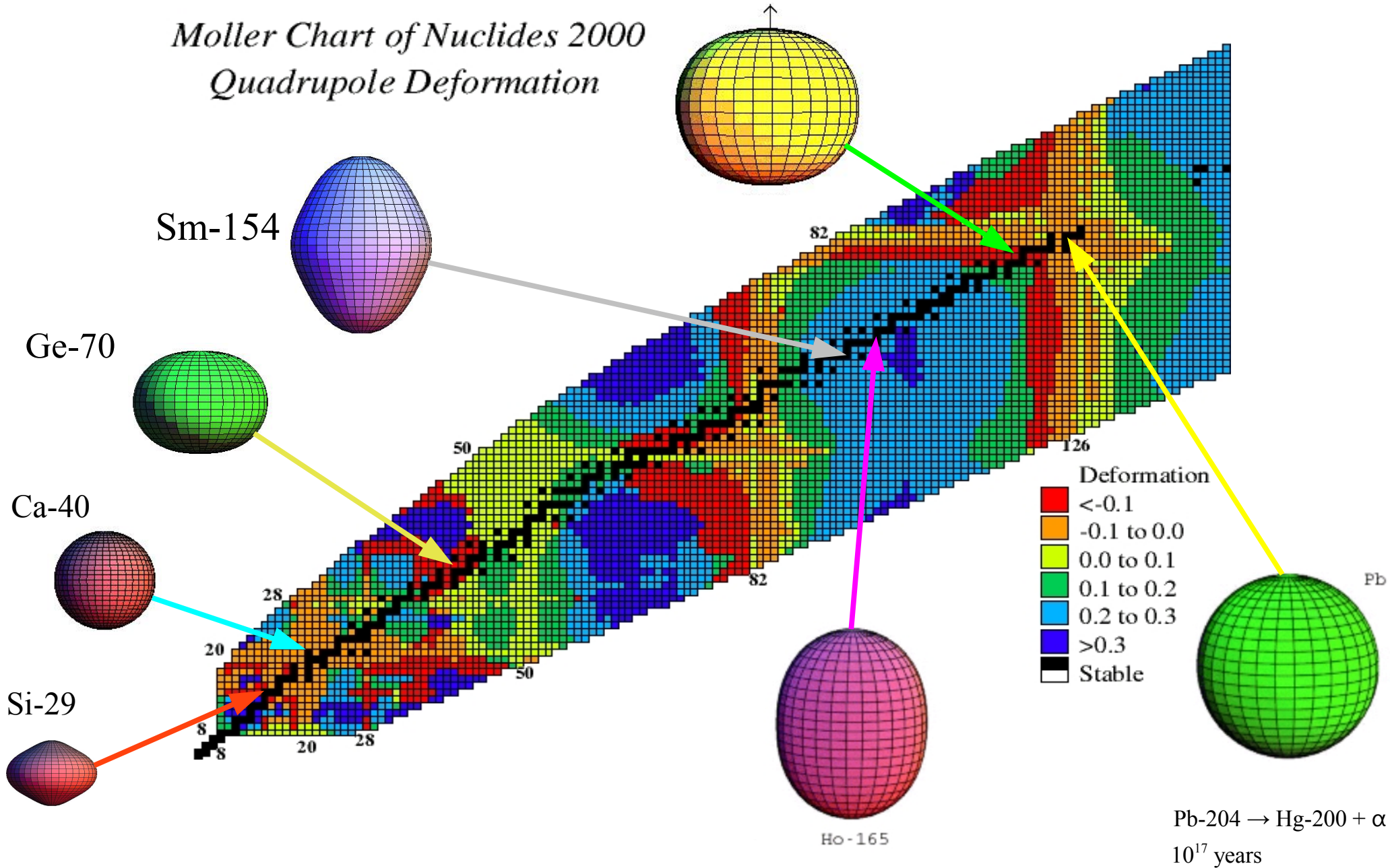
Sm-154

OUTLINE:

- Introduction to Nuclear Deformation
 - many interesting pieces, oblate, prolate...
- Collisions of deformed nuclei
 - eccentricity, elliptic flow, fluctuations....
- Conclusions...

Quadrupole deformation: Theoretical Calculation

Moller Chart of Nuclides 2000
Quadrupole Deformation



Woods-Saxon Density.

$$\rho_w(x, y, z) = \frac{\rho_0}{1 + e^{(r - R_0(1 + \beta_2 Y_{20} + \beta_4 Y_{40})) / a}}$$

- Deformation parameters:

quadrupole: $\beta_2 \rightarrow [3\cos^2(\theta) - 1] \approx Y_{20}$

octupole: $\beta_3 \rightarrow [5\cos^3(\theta) - 3\cos(\theta)]$

higher order: $\beta_4 \rightarrow [35\cos^4(\theta) - 30\cos^2(\theta) + 3] \approx Y_{40}$

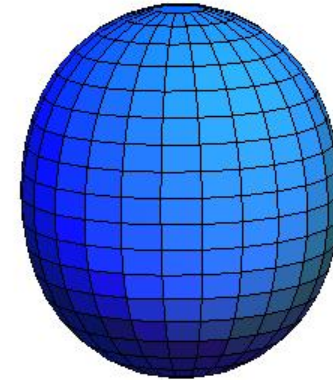
Highest order: $\beta_6 \rightarrow$ nucl-ex/0106023

- see old ref. Rev.Mod.Phys.30 pp.498-506 (1958)

Oblate/prolate shape: β_2

- $\beta_2 > 0 \rightarrow$ rugby-ball (prolate) shape.

Ne-20, Cu-63, Sm, W, U..



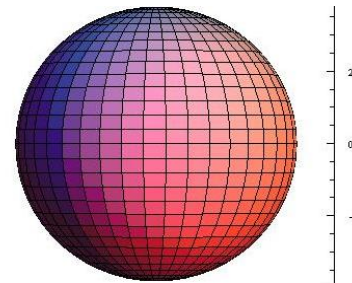
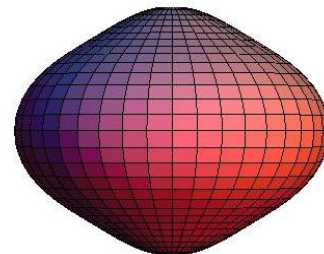
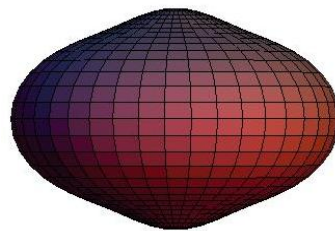
Cu-63

- $\beta_2 < 0 \rightarrow$ oblate (squeezed) shape: Si, As, Ge, Au

Si-28

Si-29

Si-30

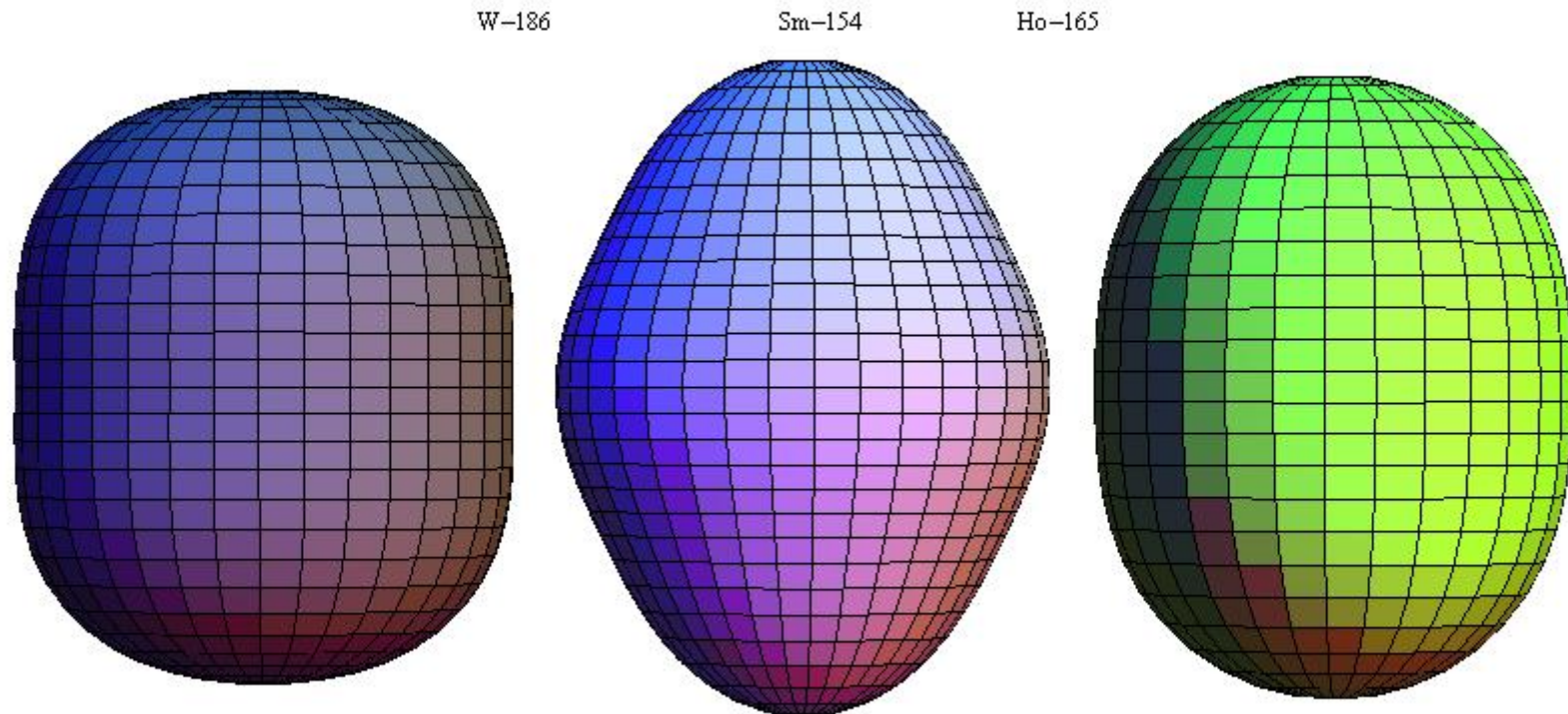


$$\beta_2 = -0.47$$

$$\beta_2 = -0.3$$

$$\beta_2 = 0.0$$

Higher-order deformation: β_4



W-186

Sm-154

Ho-165

$$\beta_4 = -0.1$$

$$\beta_4 = +0.1$$

$$\beta_4 = 0$$

Octupole deformation: β_3

- Pear-shaped deformation
→ under investigation, unstable to α -decay
- Two candidates: Sm-149 and Rn-222

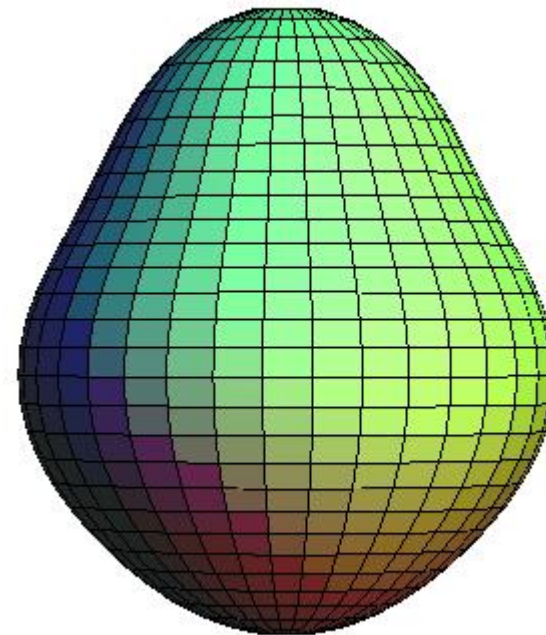
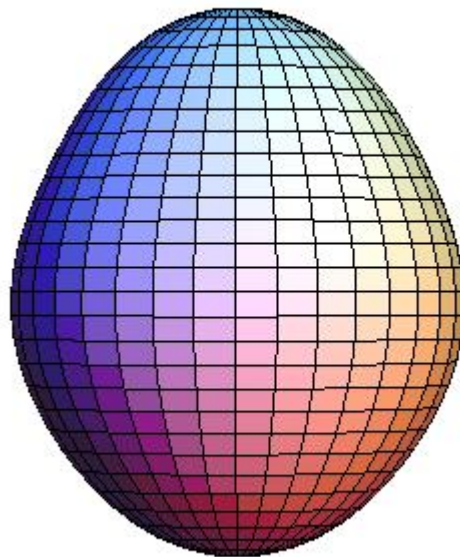
$$(\beta_3 = -0.05)$$

Sm-149

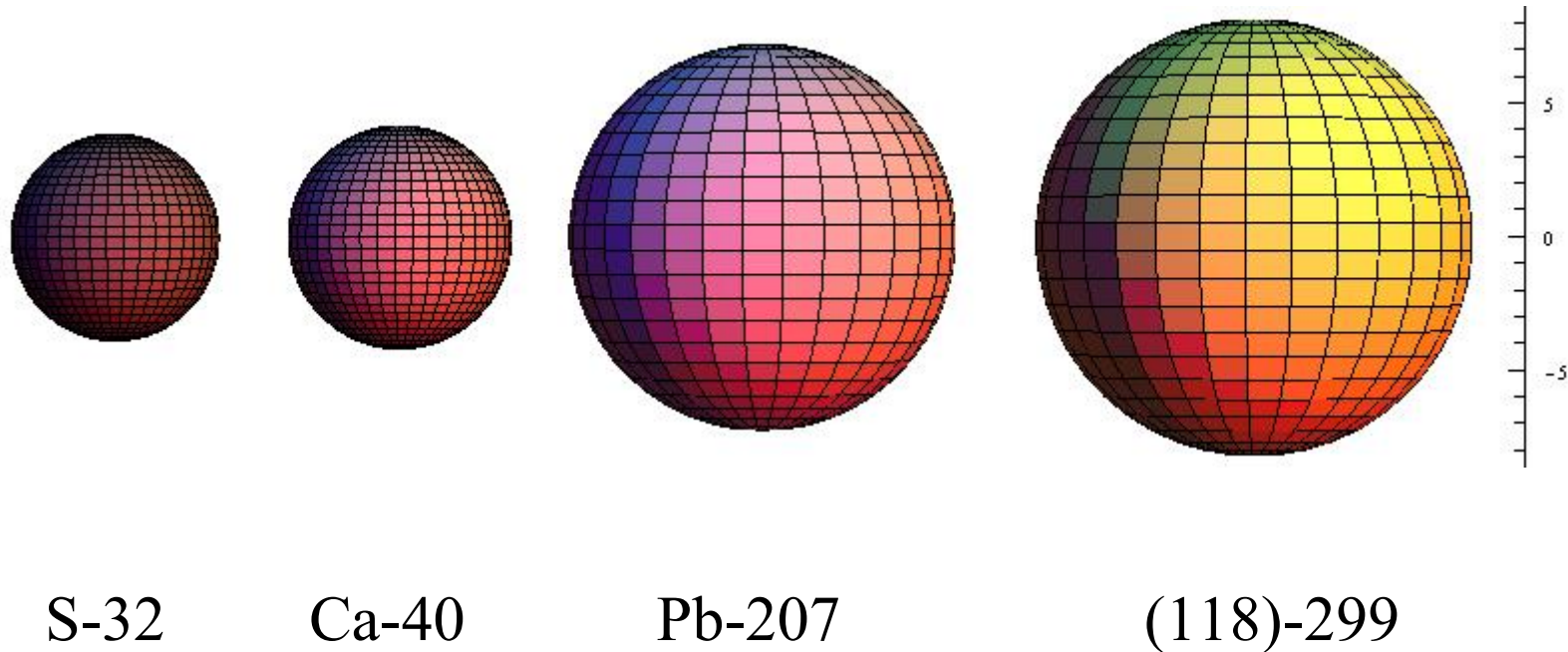
Octupole Def.

$$(\beta_3 = -0.13)$$

Rn-222



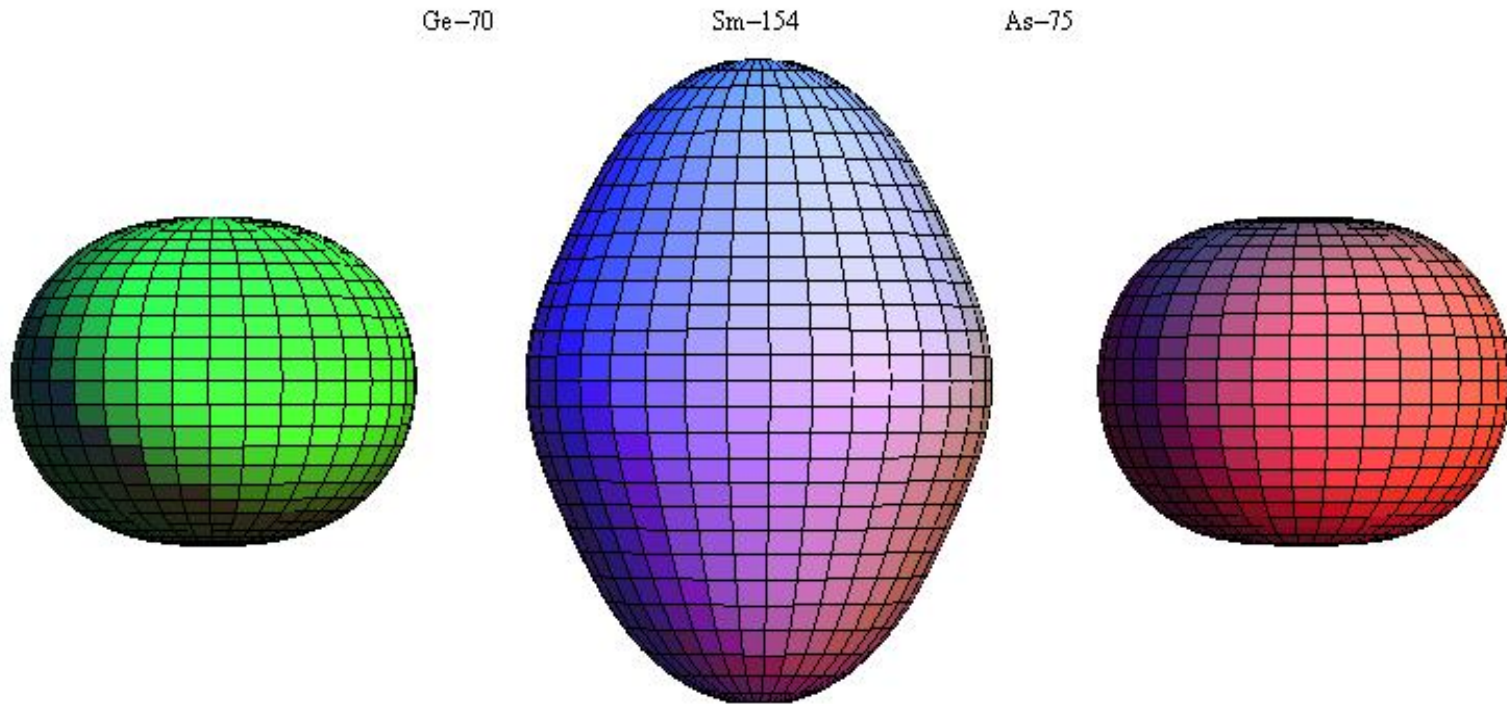
Size comparison.



Spherical nuclei = closed shells of nucleon orbitals (magic numbers).

Radius increases as $A^{1/3}$ [assuming constant baryon density].

Shape comparison I.



Ge-70

Sm-154

As-75

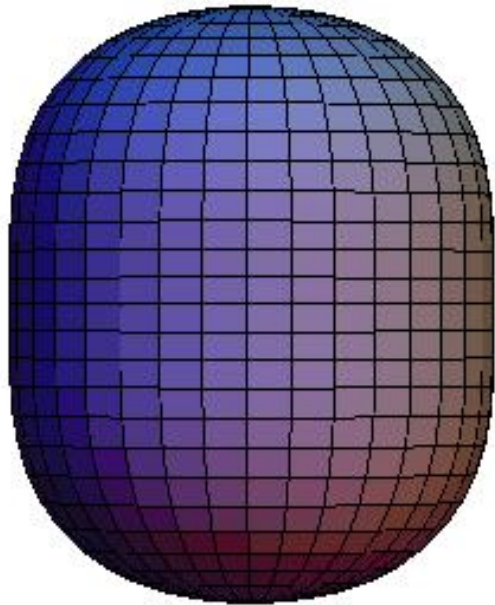
Ge-70

Sm-154

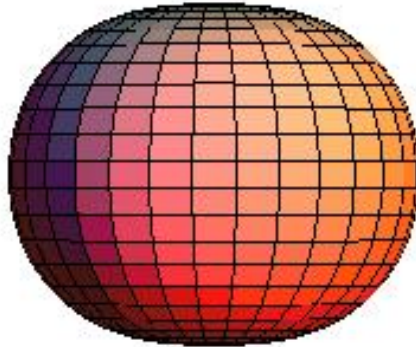
As-75

Shape comparison II.

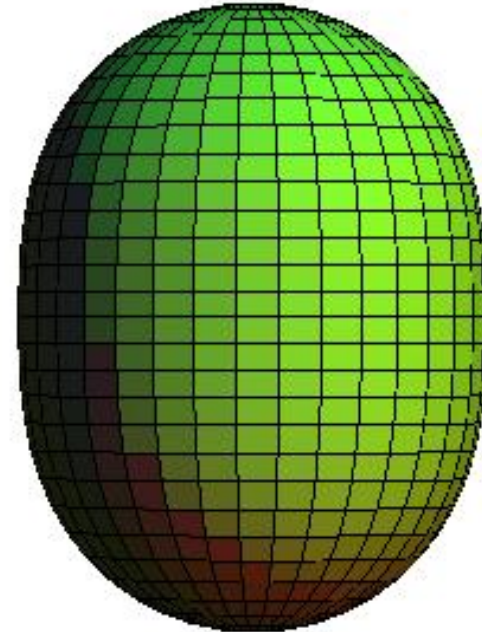
W-186



Ga-71



Tm-169

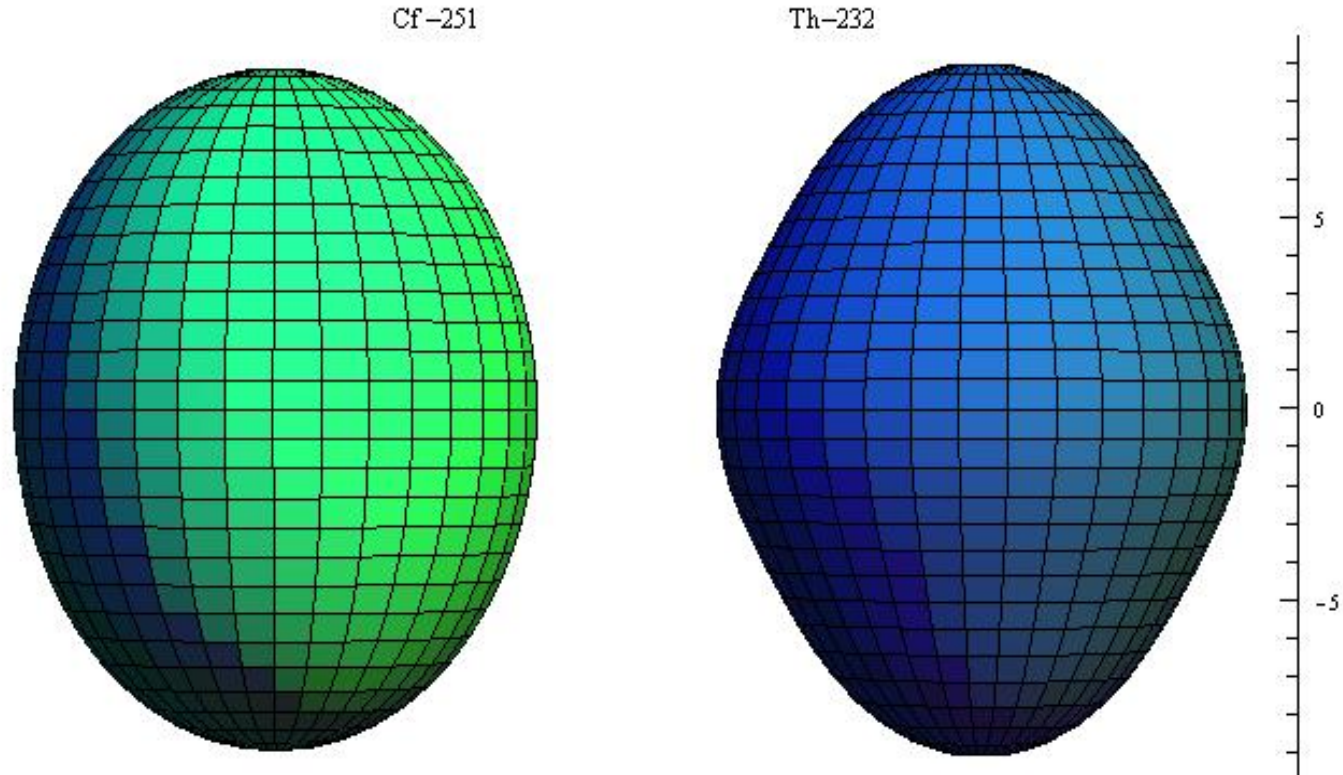


W-186

Ga-71

Tm-169

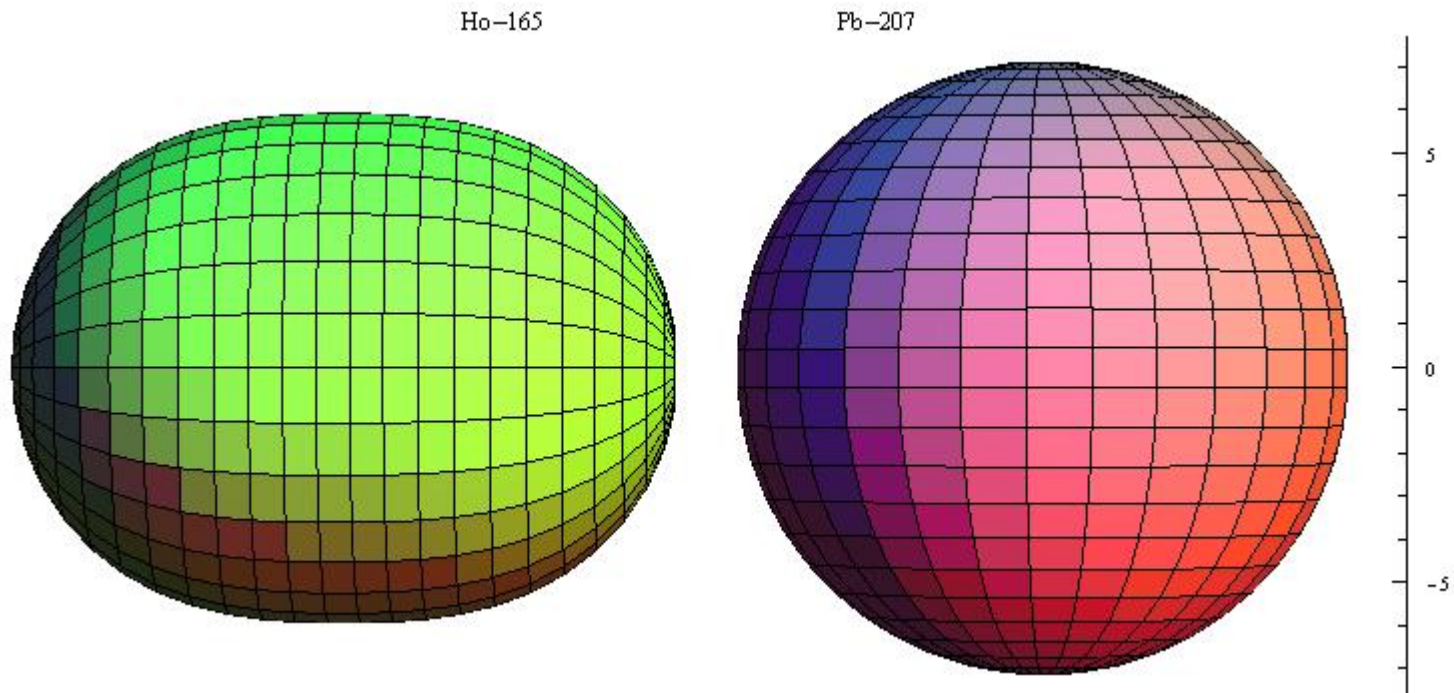
Other Pictures I.



Cf-251 and Th-232

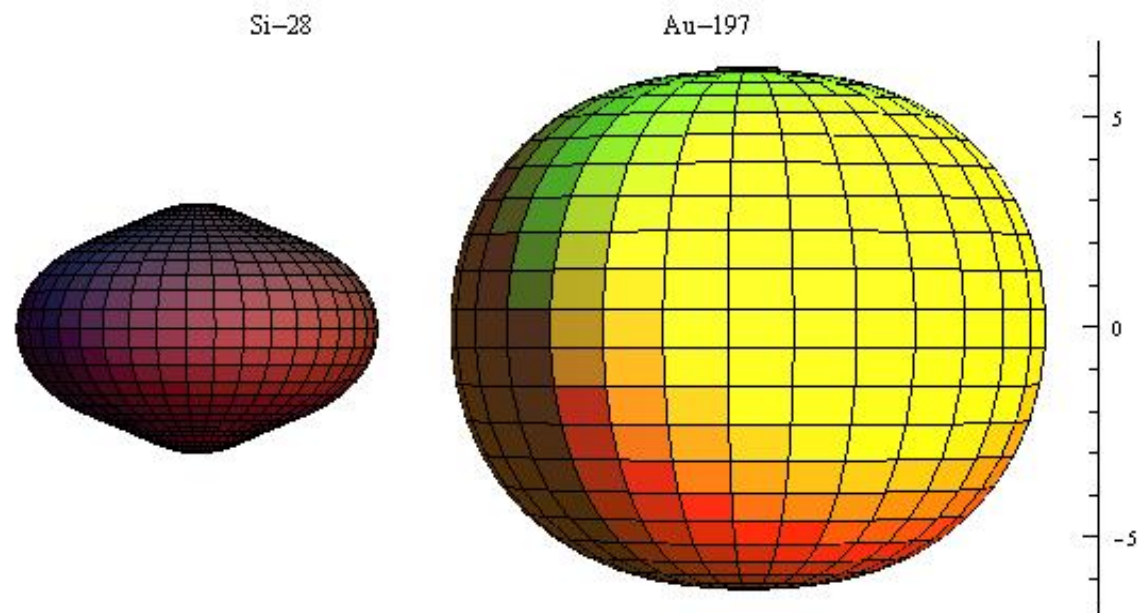
[Cf-251 in RHIC = 89800 Years half-life at 100GeV/n; for 10^{13} / beam \approx 1 decay/6s]

Other Pictures II.



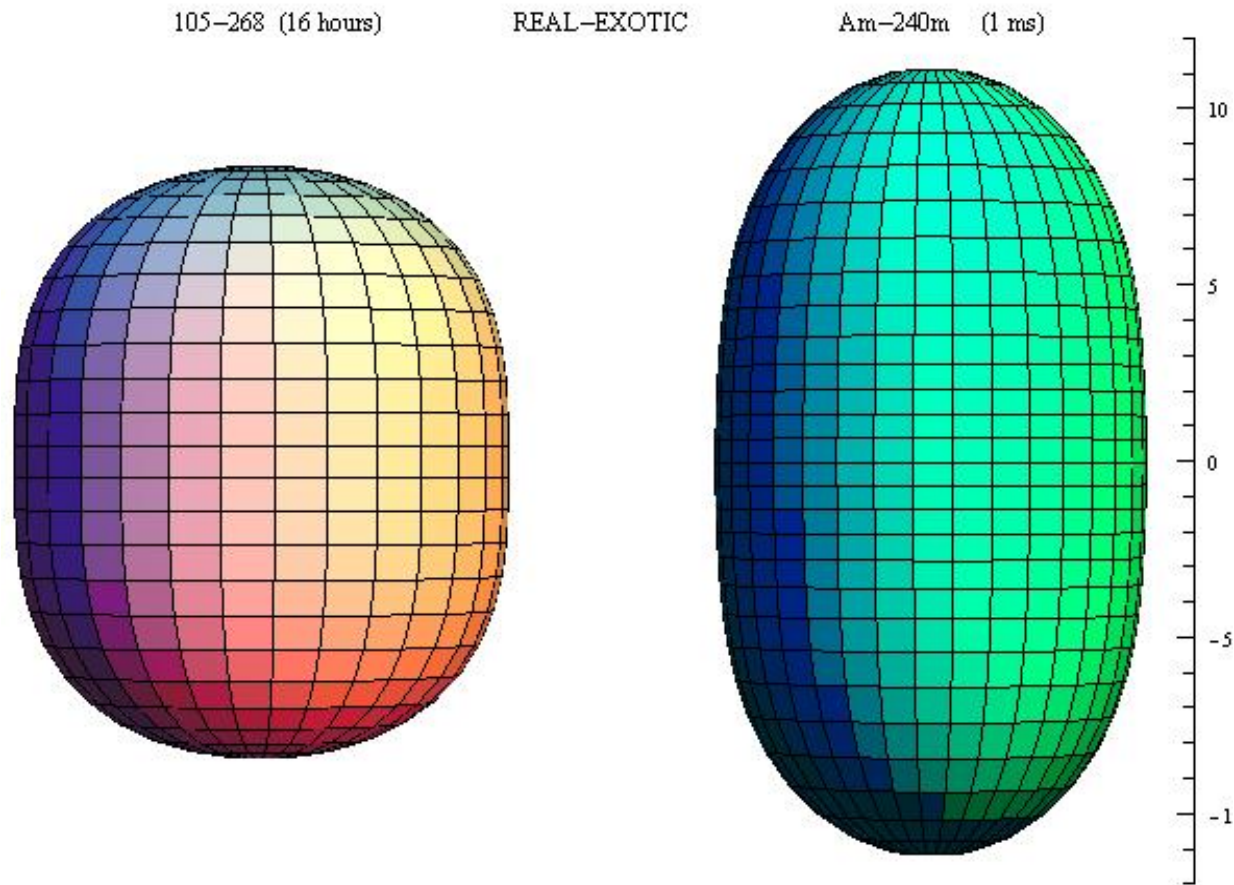
Ho-165 \rightarrow \leftarrow Pb-207
(long-polarized)

Other Pictures III.



Si-28 \rightarrow \leftarrow Au-197
(AGS)

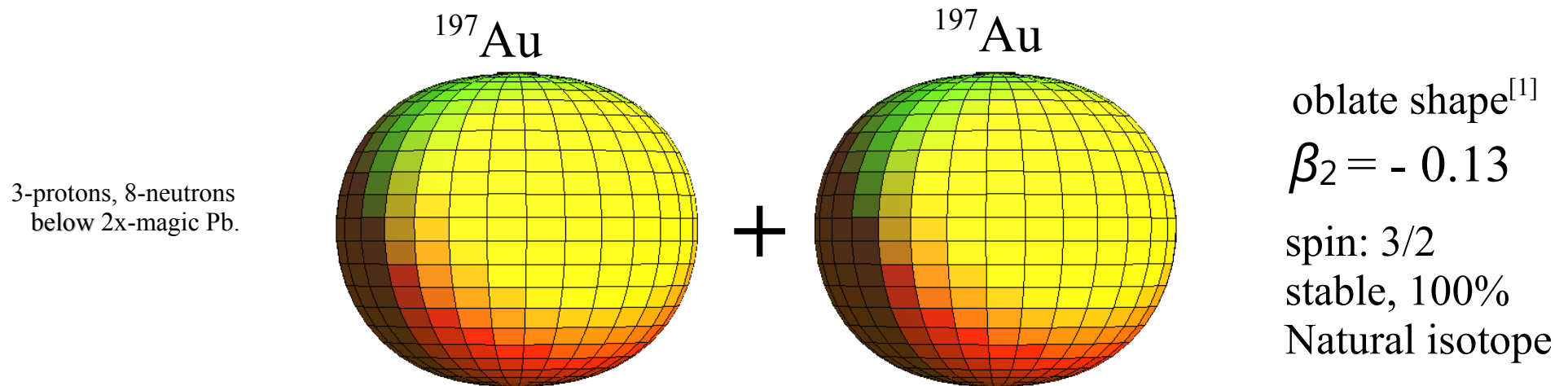
EXOTICS.



Dubnium(105)-268 (16hours)

Americium(95)-240m (1ms)

Reality: RHIC Au+Au 200GeV/n



Au-197 nucleus Deformation ?

Yes: has quadrupole moment $Q=0.59$ barn [Phys.Rev.A73(2006)022510]

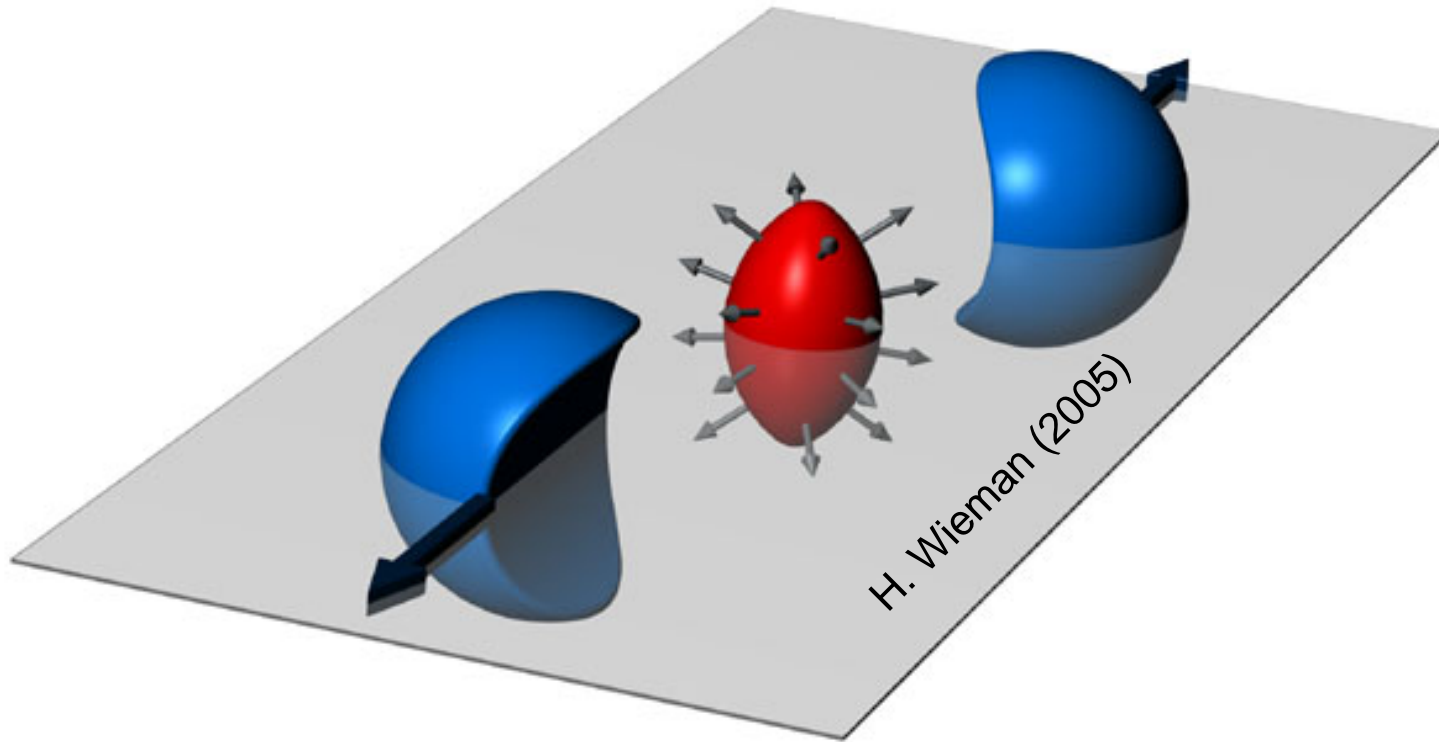
Prediction β_2 : [1] P.Moller et al. At. Data Nucl. Data Tables 59 (1995) 185.

Experiment: [2] C.Nair et al. (Giant Dipole Resonance ^{197}Au) arXiv:08114746.

Other deformed nuclei at RHIC ? **Yes:** Cu-63 & U-238.

Elliptic flow at RHIC energies:

→ origin: the initial **spatial** asymmetry.



For deformed nuclei → initial **eccentricity** is affected !!!

→ elliptic flow is affected.

Optical Glauber Model*

- Using Deformed Woods-Saxon density:

$$\rho_w(x, y, z) = \frac{\rho_o}{1 + e^{(r - R_o(1 + \beta_2 Y_{20} + \beta_4 Y_{40}))/a}}$$

- Projections $[\theta, \phi]$ in transversal plane \rightarrow
- From the overlap of colliding nuclei:

- Baryon density

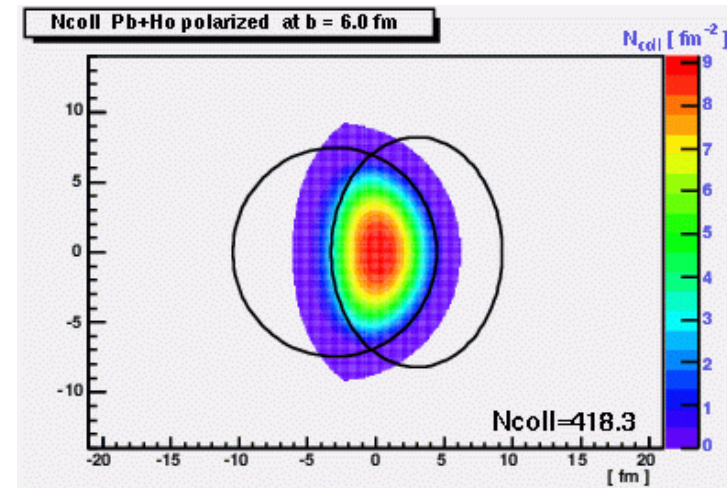
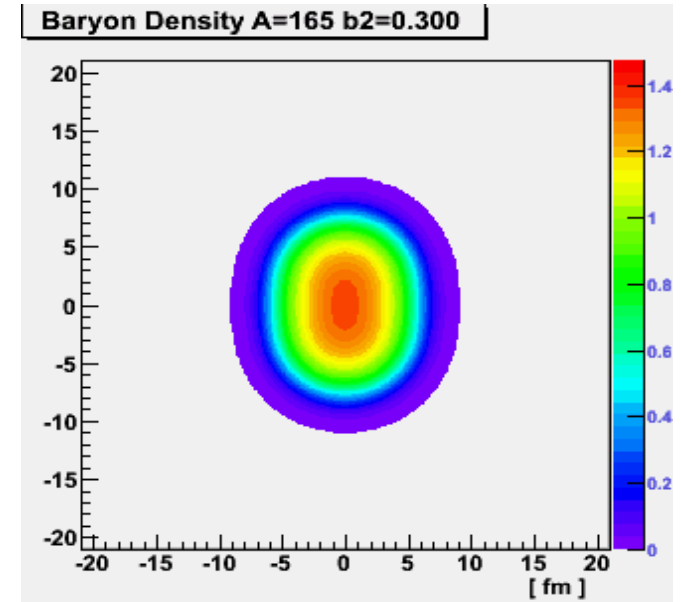
- N_{part} : participant density $\rho_{\text{part}}(x, y)$

- N_{coll} : $\rho_{\text{coll}}(x, y)$ binary collisions density

- Obtain eccentricity $\varepsilon = \frac{\sqrt{(\sigma_y^2 - \sigma_x^2)^2 + 4\sigma_{xy}^2}}{\sigma_y^2 + \sigma_x^2}$

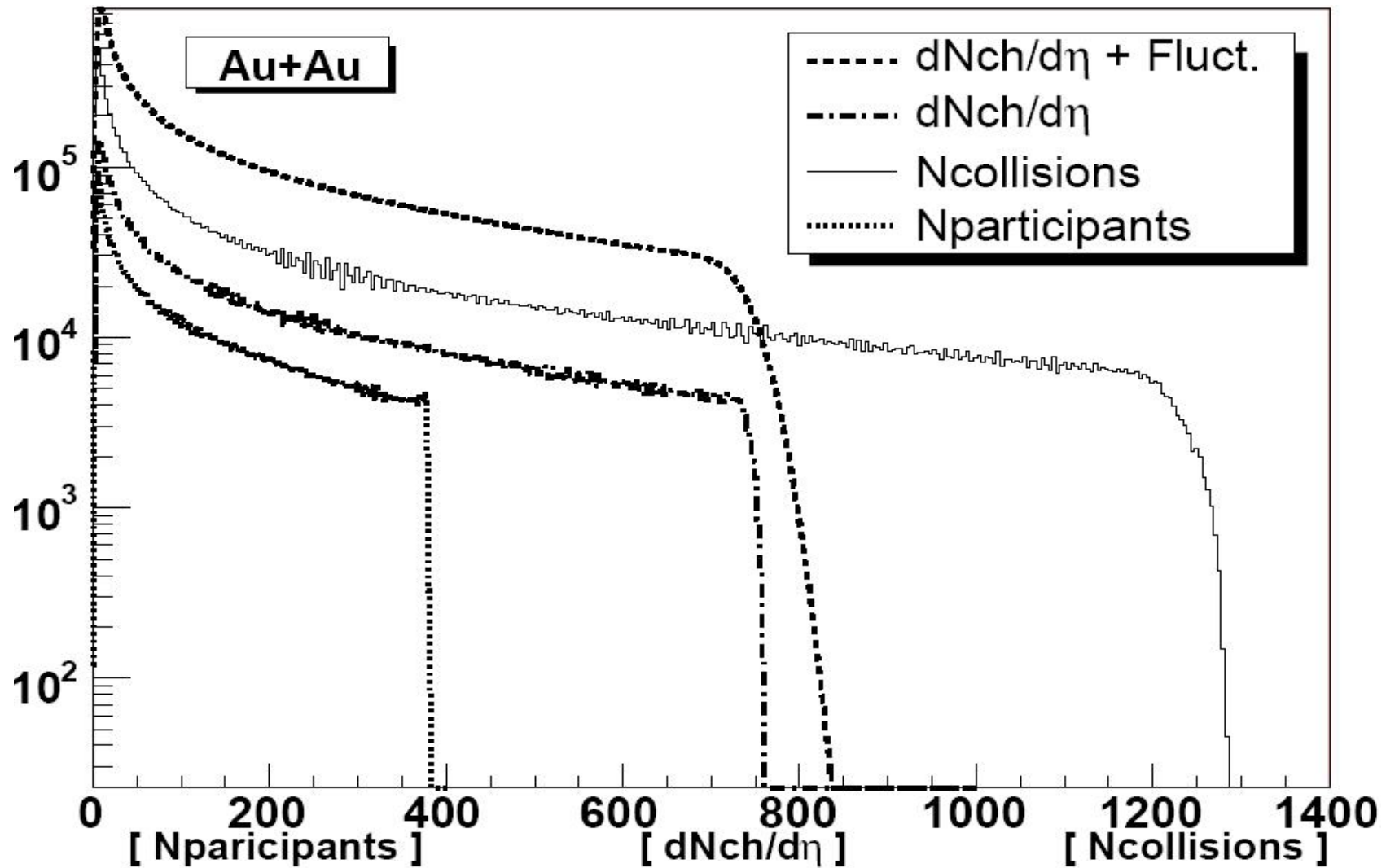
- Multiplicity:

$$dN_{\text{ch}}/d\eta = (1 - x) \cdot n_{pp} \frac{N_{\text{part}}}{2} + x \cdot n_{pp} N_{\text{coll}}$$



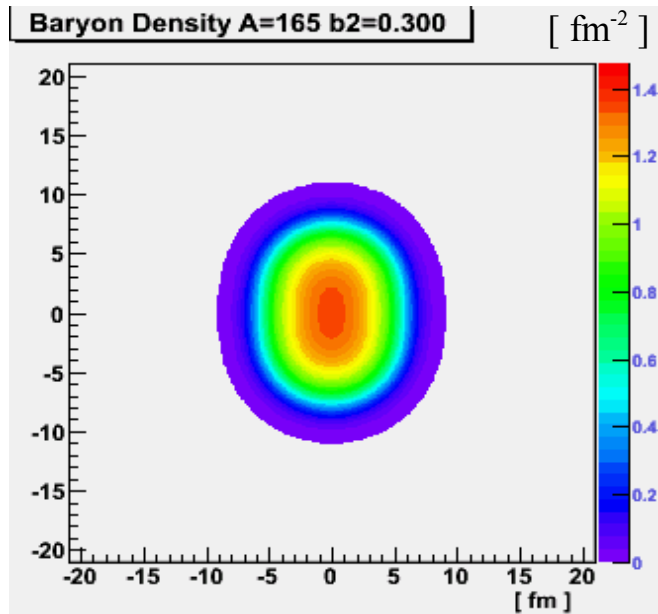
* Phys. of Atom. Nucl. 71 (2008) 1609

Distributions N_{part} , N_{coll} , N_{ch} from Opt.GM



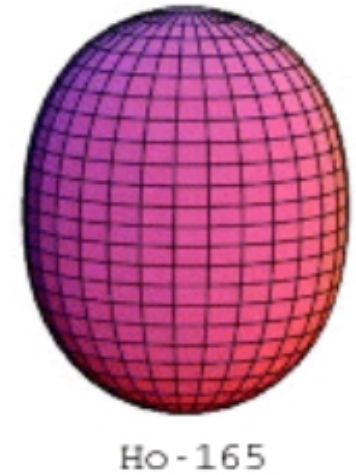
Two-component $dN_{\text{ch}}/d\eta$: Phys.Lett.B507(2001)121; $n_{pp} = 2.29$ and $x = 0.13$

Eccentricity in collisions of prolate nuclei.

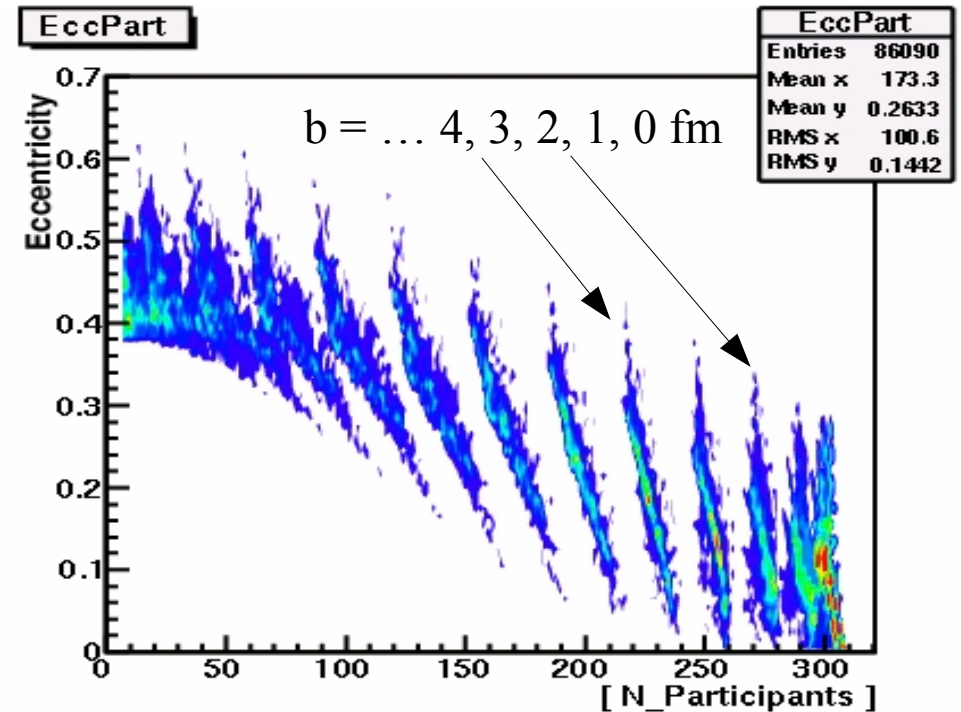
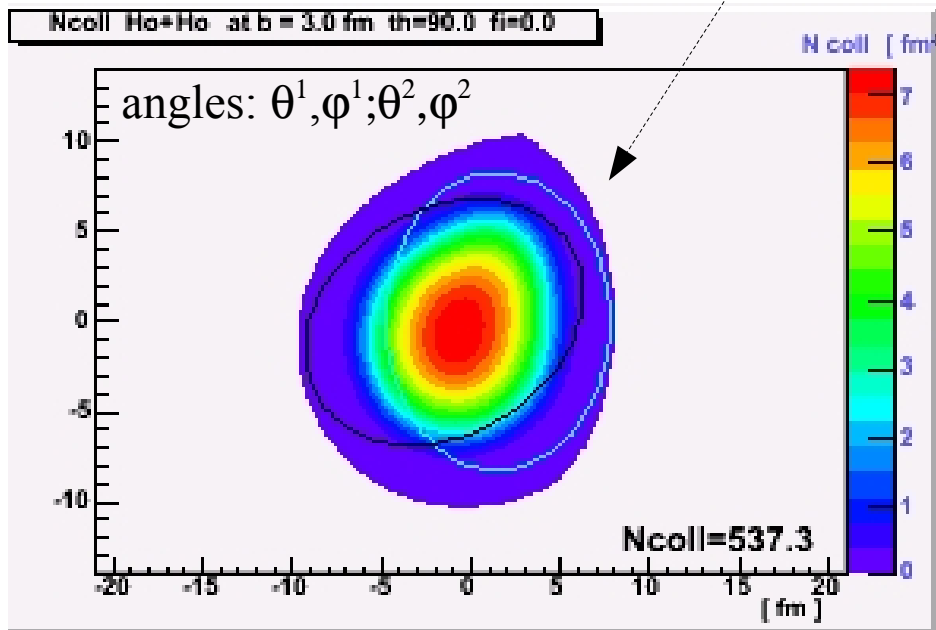


- Ho-165 ($\beta_2 = +0.3$)

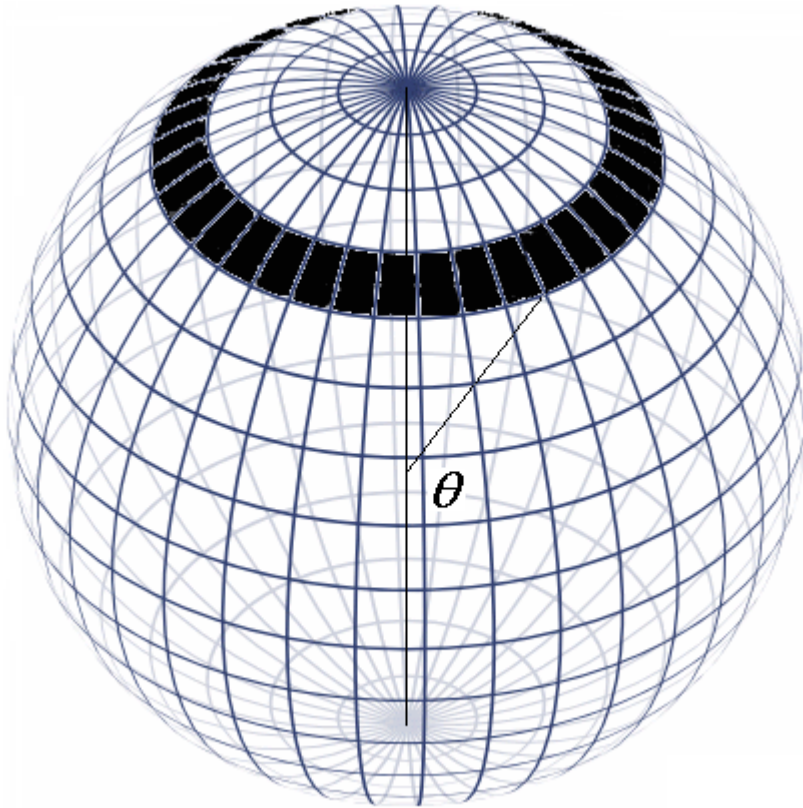
$$\varepsilon_{\text{part}} = \frac{\sqrt{(\sigma_y^2 - \sigma_x^2)^2 + 4\sigma_{xy}^2}}{\sigma_y^2 + \sigma_x^2}$$



- Pear-shaped overlaps !
- Fluctuations $\varepsilon[\theta^1, \varphi^1; \theta^2, \varphi^2]$ at given [b]



Technical detail: Random orientation of nucleus.



Random orientation:

= random distribution of points where main axis (spin) crosses the surface of the sphere.

Probability is proportional to the area.

Area dS corresponding to $d\theta d\varphi$ is:

$$dS = R \sin(\theta) d\theta d\varphi$$

→ $P(\theta) = \sin(\theta)/2$ (normalized to 1.)

→ $P(\varphi) = \text{const.}$ (random φ angle)

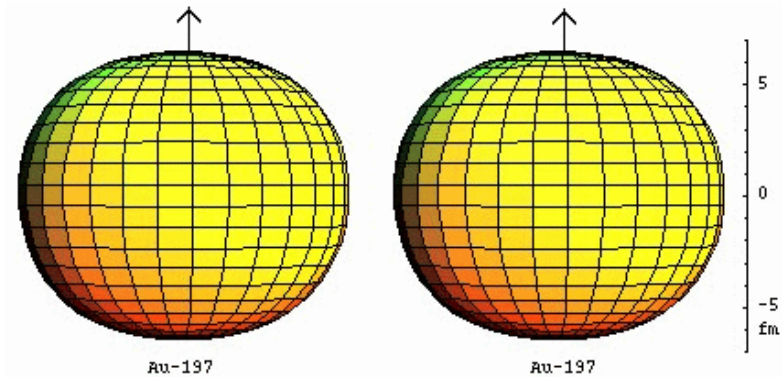
→ angle θ is not random.

Random orientation means random φ , and $\sin(\theta)$ distributed θ angle.

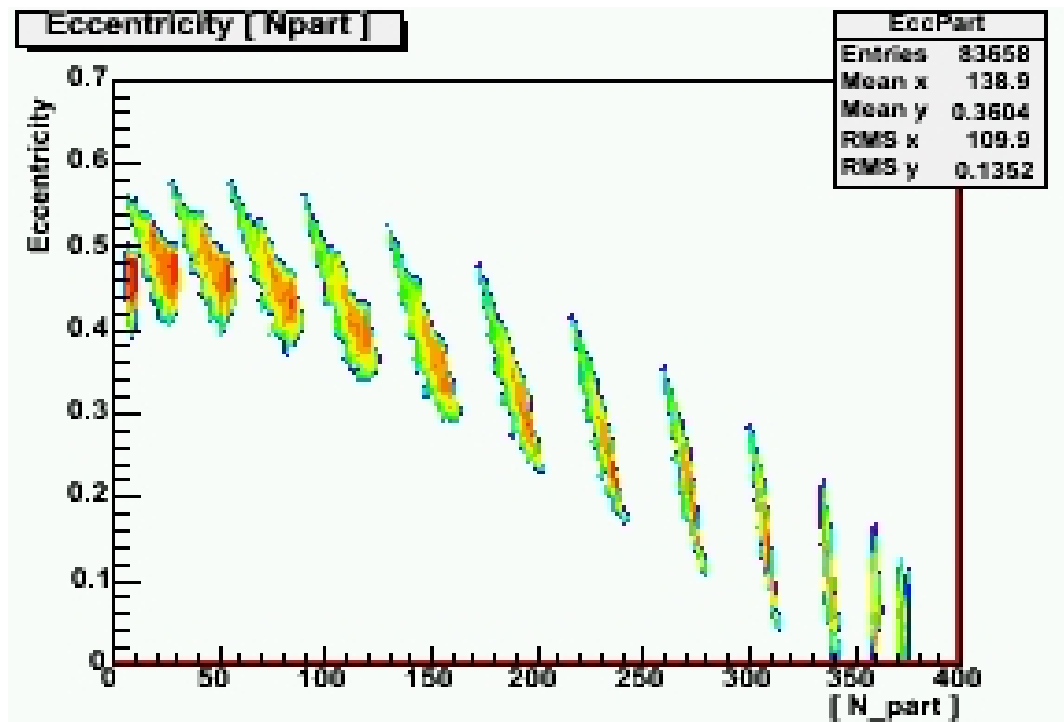
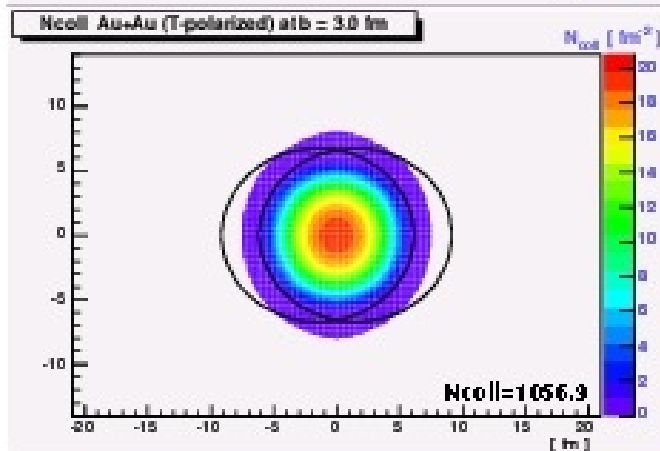
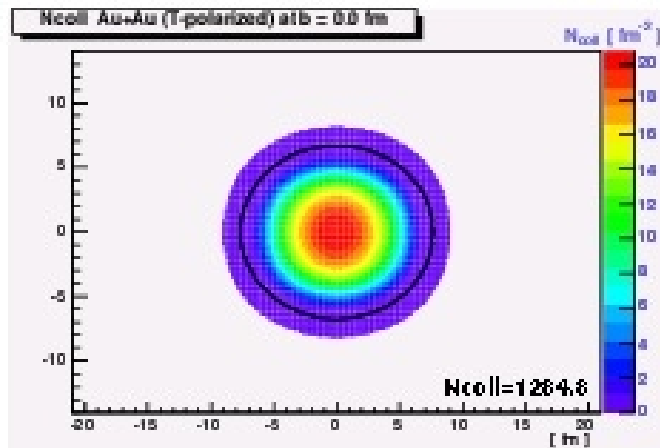
Eccentricity in collisions of oblate nuclei.

- Au-197 (predicted $\beta_2 = -0.13$)

Zero eccentricity at $b=3\text{fm}$
and non-zero ϵ for $b=0\text{fm}$



Eccentricity fluctuates again !



OBSERVATION:

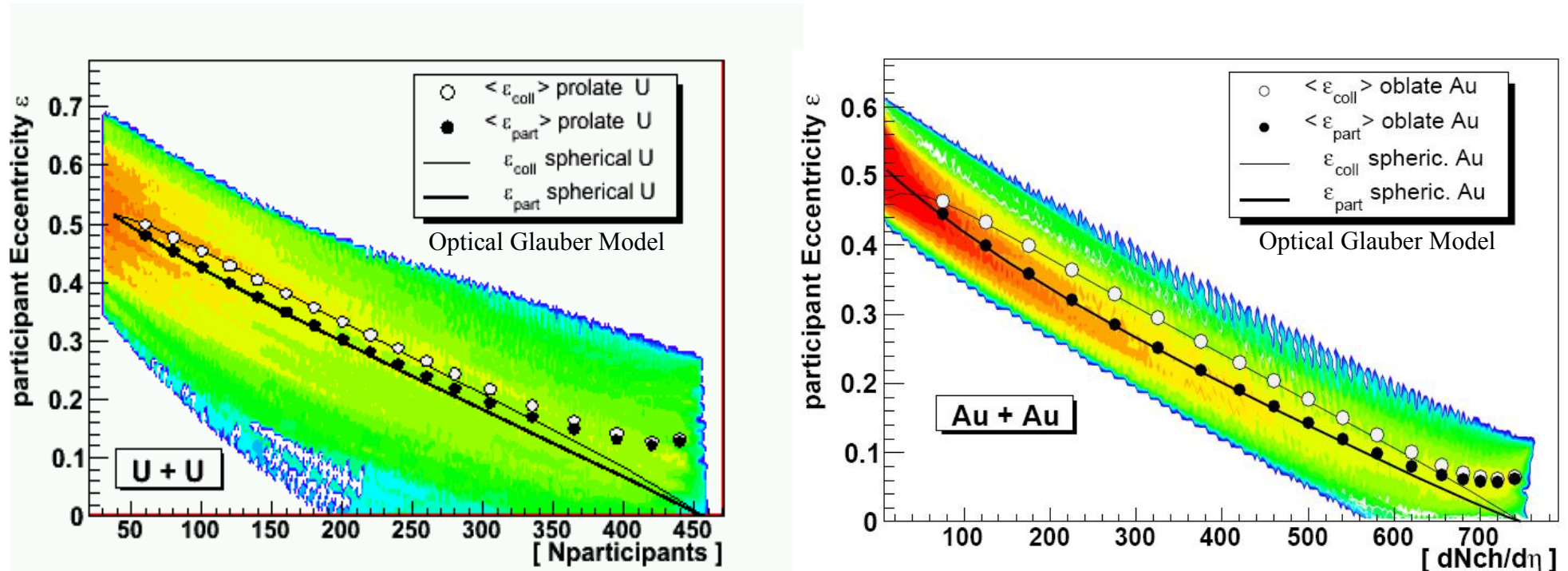
Deformation of nuclear shape

increases

$\varepsilon \rightarrow v_2 =$ Elliptic flow **fluctuations**.

(at given fixed collision centrality).

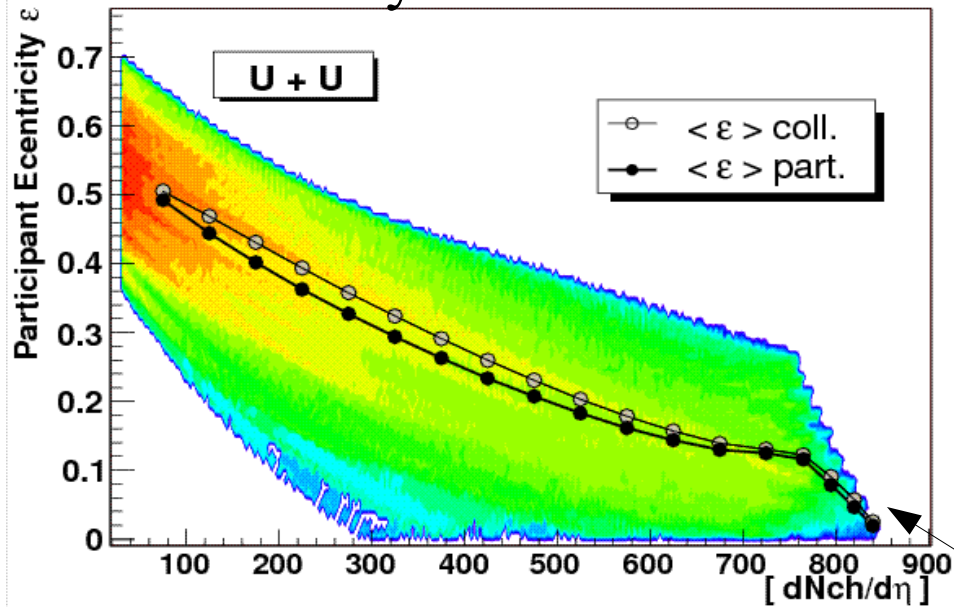
What happens with $\langle \varepsilon \rangle$ due to deformation in UU & AuAu ?



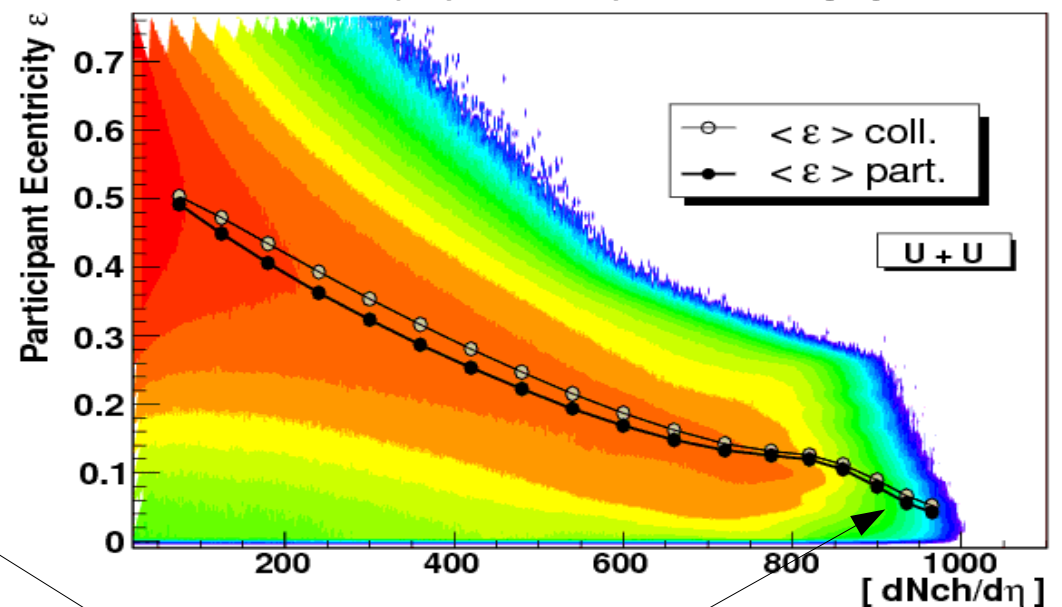
- In noncentral collisions $\langle \varepsilon \rangle$ stays unchanged
 → central coll: Increased $\langle \varepsilon \rangle$ due to deformation.
 + additionally, deformation increases eccentricity fluctuations.

UU collisions from Opt.GM: $dN_{ch}/d\eta$

Only deformation



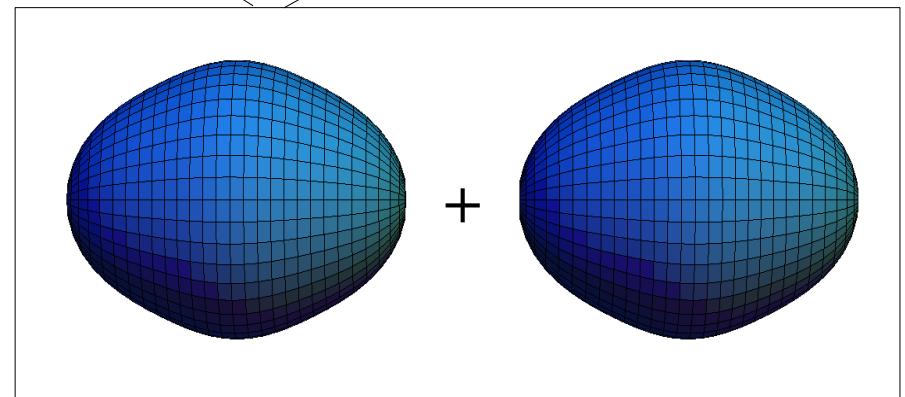
Deformation + FLUCT



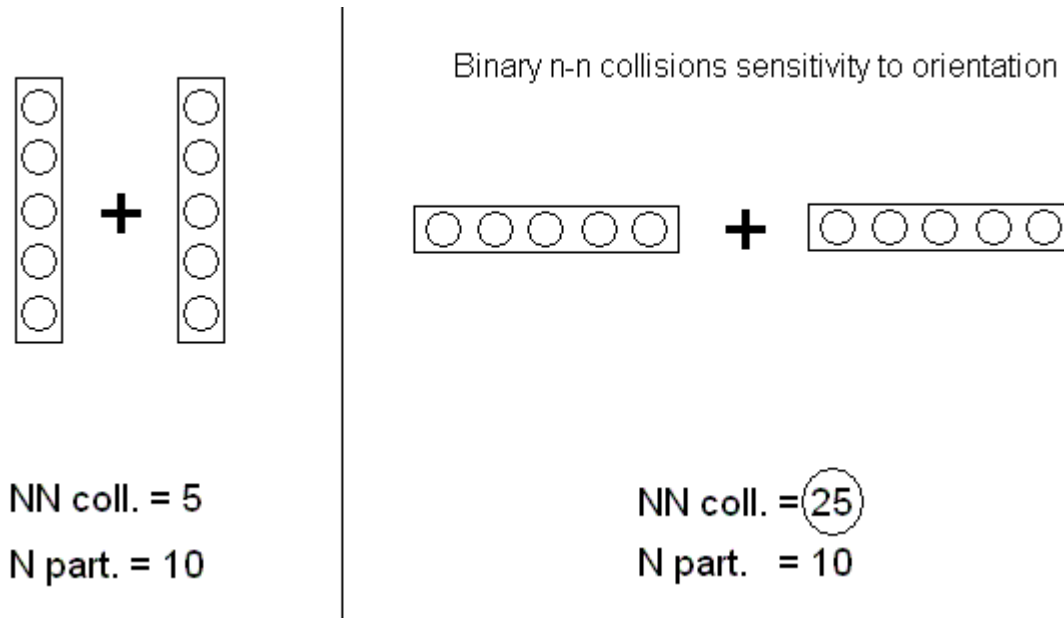
→ Cusp in $\langle \varepsilon \rangle$ for very-central collisions (large $dN_{ch}/d\eta$).

Highest multiplicity: $dN_{ch}/d\eta$
is observed
for longitudinally polarized,
central $b=0$ fm collisions.

→ eccentricity cusp ←



Why $(dN_{ch}/d\eta)$ sensitive to orientation ?



→ $dN_{ch}/d\eta$ depends on orientation due to NN_{coll}

$$dN_{ch}/d\eta = (1 - x) \cdot n_{pp} \frac{N_{part}}{2} + x \cdot n_{pp} N_{coll}$$

↙

N_{part} is **not** sensitive to orientation v_2 [N_{part}] not interesting.

→ Study: v_2 [$dN_{ch}/d\eta$] (in central collisions) !

Glauber Monte Carlo Simulation*

[*] QM09 poster (LBNL + JINR)

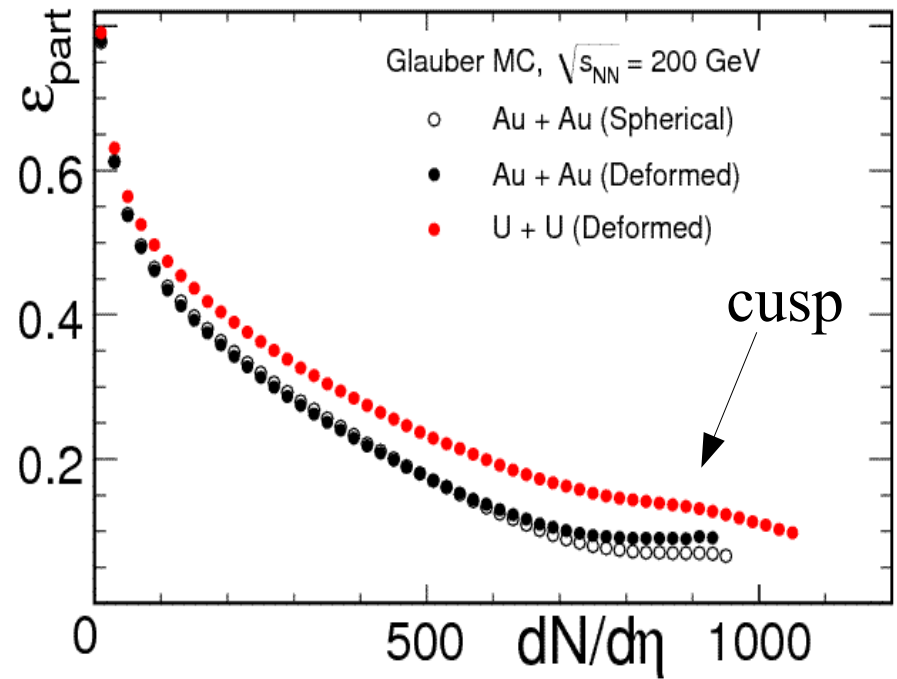
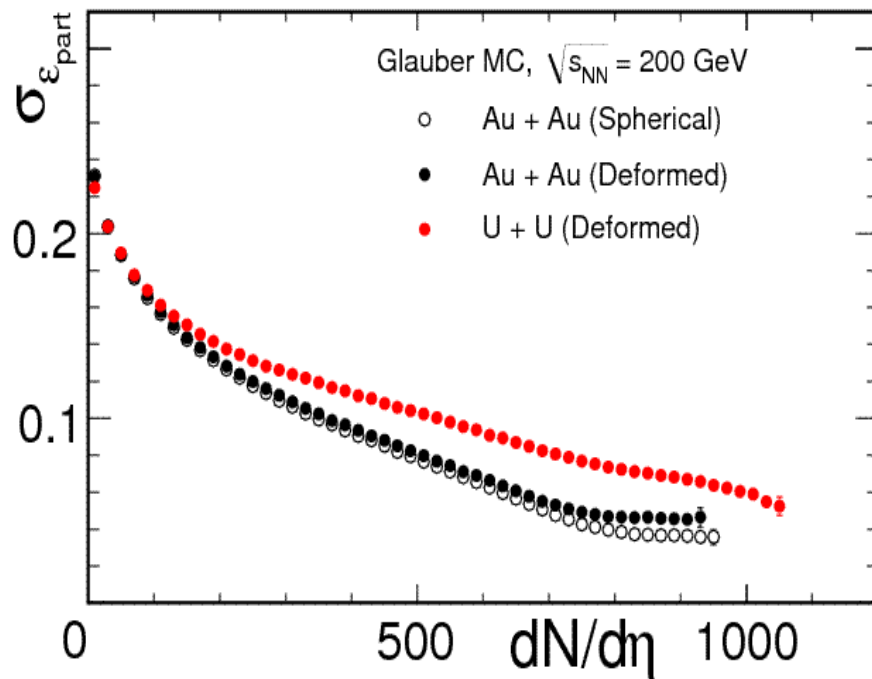
Eccentricity fluctuations: $\sigma_\epsilon = \sqrt{\sigma_{\beta_2}^2 + \tilde{\sigma}_\epsilon^2}$

→ finite number of interacting nucleons: $\tilde{\sigma}_\epsilon$

→ ground-state deformation of coll. nuclei: σ_{β_2}

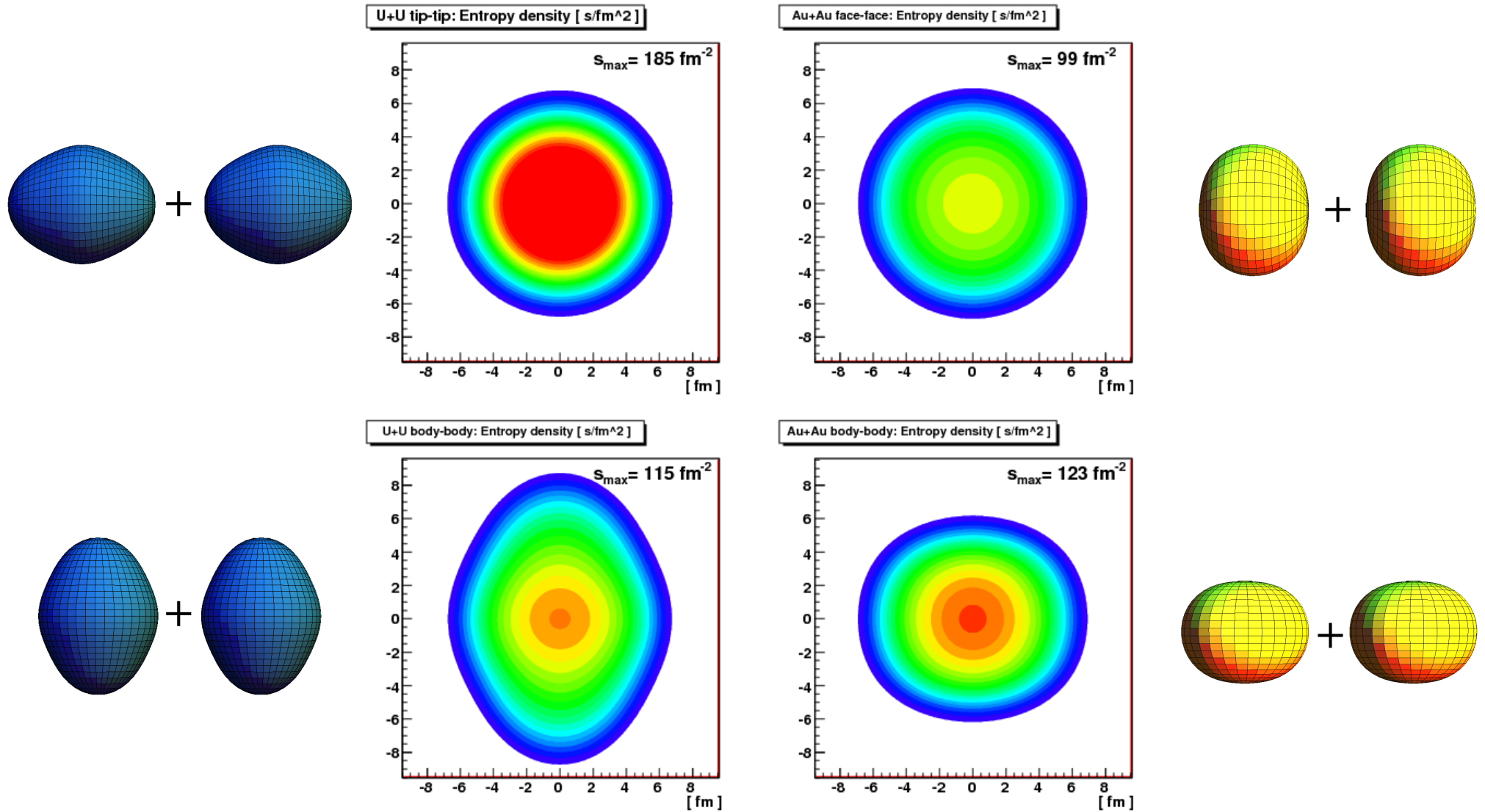


Quark Matter 2009



Effects predicted by Optical Glauber Model → confirmed.

Comparison of Entropy Density in Au+Au vs. U+U

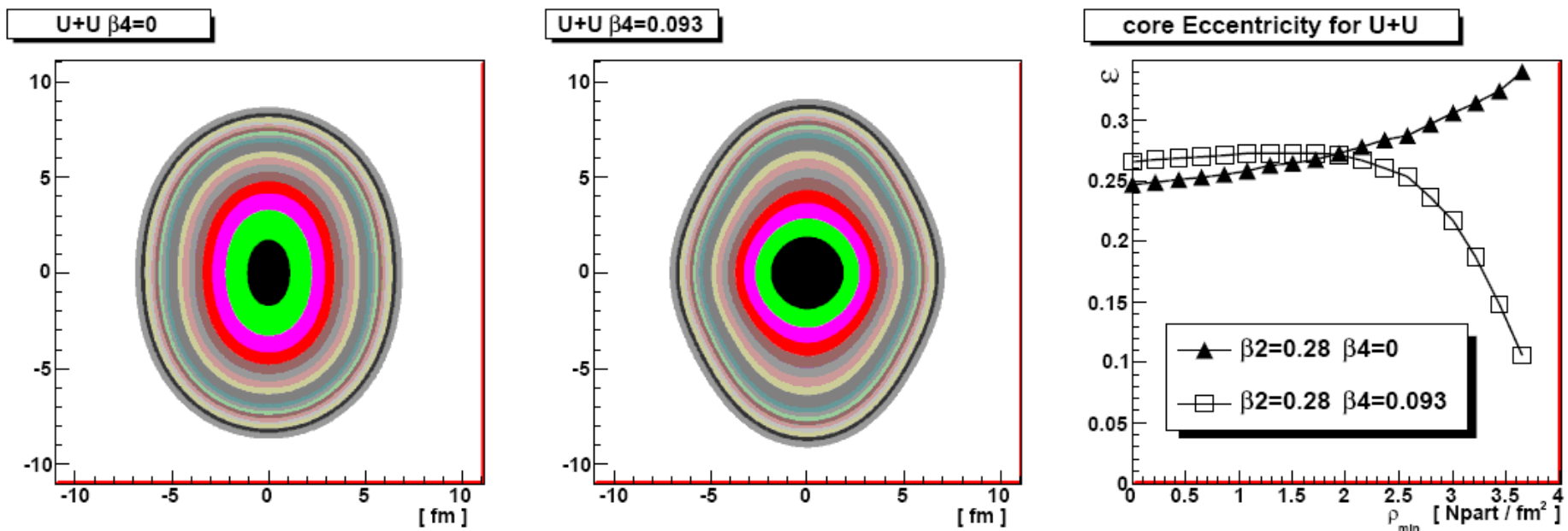


Entropy density: $\rho_s(x, y) = \kappa_s [\alpha \rho_{\text{part}}(x, y) + (1 - \alpha) \rho_{\text{coll}}(x, y)]$

β_4 deformation Parameter for ^{238}U

- Moller et al. [1] prediction for ^{238}U is: $\beta_4 = + 0.093$
 - spatial distribution of nucleons is modified
 - participant eccentricity is modified
 - final v_2 strength can be modified ! [hydro calc. needed].

U+U participant density in transversal plane [fm^{-2}].



[1] Moller et al. At. Data Nucl. Data Tables 59, 185 (1995)

Polarization of ^{238}U beam ? No way.

Spin of $^{238}\text{U} = 0^+$

Magnetic moment $\mu = 0$

Quadrupole moment: unknown.

→ consider ^{238}U beam unpolarized.

CONCLUSIONS:

- **Nuclei collided at AGS/SPS/RHIC are deformed:**
→ Si-28, Cu-63, In-115, Au-197, (U).
- **Elliptic flow is affected in deformed nuclei collisions:**
→ eccentricity: $\langle \varepsilon \rangle$ increased in Au+Au central
+ cusp for Ho+Ho and U+U predicted
→ fluctuations: $\sigma(\varepsilon)$ increased !
- β_4 deform. parameter for ^{238}U is important !
- **Study of deformation effects is needed to understand properties of partonic matter created at RHIC.**